

Going beyond Pentacene: Photoemission Tomography of a Heptacene Monolayer on Ag(110)

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KEYWORDS

Surface science, acenes, organic semiconductors, angle-resolved photoemission spectroscopy (ARPES), photoemission tomography (PT), ultrathin films.

ABSTRACT

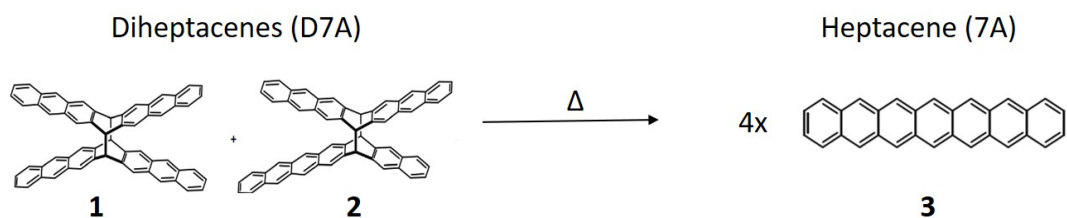
Longer acenes such as heptacene are promising candidates for optoelectronic applications but are unstable in their bulk structure as they tend to dimerize. This makes the growth of well-defined monolayers and films problematic. In this article, we report the successful preparation of a highly oriented monolayer of heptacene on Ag(110) by thermal cycloreversion of diheptacenes. In a combined effort of angle-resolved photoemission spectroscopy and density functional theory calculations, we characterize the electronic and structural properties of the molecule on the surface in detail. Our investigations allow us to unambiguously confirm the successful fabrication of a highly oriented complete monolayer of heptacene and to describe its electronic structure. By comparing experimental momentum maps of photoemission from frontier orbitals of heptacene and pentacene, we shed light on differences between these two acenes regarding their molecular orientation and energy level alignment on the metal surfaces.

INTRODUCTION

π -conjugated molecules have gained a crucial role in the development of organic electronic technologies. Among compounds which have attracted ongoing interest is the class of acenes, i.e., polycyclic aromatic hydrocarbons consisting of n linearly fused benzene rings, which we are going to notate as nA in the following. Here, the commercially available tetracene (4A, $n = 4$) and pentacene (5A, $n = 5$) are the most prominent members, which have been studied extensively over the years.^{1, 2} Longer acenes are attractive due to their narrow band gap between the highest occupied and lowest unoccupied molecular orbitals (HOMO and LUMO, respectively),³ low reorganization energy⁴ and high charge carrier mobility.⁵ These optical and electronic properties strongly depend on the number of benzene units and can thus be tuned by the number of fused rings.^{6,7} More importantly, longer acenes may serve as the atomically precise model systems for the narrowest possible zigzag-edge graphene nanoribbons with potential applications in spintronics and plasmonics.^{8,9} However, the experimental study and application of acenes longer than 5A has so far been limited by their poor solubility and their pronounced tendency towards degradation,¹⁰⁻¹² making synthesis and handling under conventional conditions difficult. Carried out within a solid matrix, the production of acenes from hexacene upwards has been demonstrated.¹³⁻¹⁵ In the last decade, on-surface synthesis has been established as an alternative approach towards the generation of compounds hard to access otherwise.¹⁶⁻¹⁹ Linear acenes up to dodecacene²⁰⁻²⁶ have recently been synthesized in this way, but on-surface synthesis of acenes is mainly restricted to submonolayer coverages on crystalline surfaces. The largest acene, for which film formation on crystalline surfaces was achieved and adequately characterized, is hexacene.^{27, 28}

Its next higher neighbor heptacene **3** can be obtained from thermally induced cycloreversion of a D_{2h} symmetrical **1** and a C_s symmetrical **2** isomer (Scheme 1), which can be synthesized via a modified Meerwein–Ponndorf–Verley (MPV) reduction²⁹ (cf. Supporting Information (SI)). It was shown that thermolysis of the diheptacene (D7A) mixture results in sublimation of 7A, and it was suggested that this may enable the deposition of 7A films on various substrates via organic vapor deposition under ultra-high vacuum conditions.²⁹ Until now, however, heptacene (7A) monolayers on monocrystalline surfaces have never been reported.

Scheme 1. Synthesis of heptacene (7A) by thermal cycloreversion of diheptacenes (D7A). **1** and **2** denote D_{2h} symmetrical and C_s symmetrical D7A, respectively, **3** denotes 7A.



In general, applications in electronic devices rely on well-defined interfaces between organic films and inorganic substrates and on a detailed knowledge of the electronic structure at the interface,³⁰ which can be acquired by angle-resolved photoemission experiments (ARPES), in particular using photoemission tomography (PT).^{31, 32} The momentum map of photoemission, i.e., intensity of photoemission as a function of k_x , k_y at fixed binding energy (E_b), reflects the square of the Fourier transform of the molecular orbitals, from which the electrons were emitted.³³ Thus, they not only allow the emission to be assigned to particular orbitals, they also allow the identification of molecules and their orientation on the surface.³⁴ Here, we have applied PT to

confirm that indeed a well-oriented 7A monolayer can be grown on Ag(110) via thermal evaporation from D7A powder. This allows us to investigate the adsorption geometry of the prepared monolayer and orbital structure of 7A in great detail.

METHODS

Sample preparation and photoemission experiments were conducted in an ultra-high vacuum vessel (10^{-10} mbar). The Ag(110) substrate was prepared by cycles of sputtering (Ar^+ ions at 1 kV) and annealing (450 °C). Diheptacene loaded into a Knudsen type molecular evaporator decomposed at 300°C resulting in molecular heptacene evaporating in vacuum and adsorbing on Ag(110) located in front of the evaporator and kept at room temperature. The molecular flux of heptacene was calibrated using a quartz microbalance. Photoemission experiments were performed at the Metrology Light Source of the Physikalisch-Technische Bundesanstalt (Berlin, Germany) at the insertion device beamline using a toroidal electron-energy analyzer.³⁵ A photon energy of 35 eV and an incidence angle of $\chi = 40^\circ$ with respect to the surface normal were used. The polarization direction is in the specular plane, which is also the measured photoelectron trajectory. The emitted electrons were simultaneously recorded at polar angles of -80° to 80° with respect to the surface normal. Note that this direction in the shown momentum maps corresponds to the maximal molecular emission relative to the substrate emission. To measure the momentum maps for a chosen binding energy, the sample was rotated in the azimuthal direction in 2° steps, which results in a full photoelectron distribution in the (k_x, k_y) plane perpendicular to the sample normal. The energy distribution curves were obtained by integration of photoemission intensity over the entire available $k_{||}$ range.

All calculations were performed within the framework of density functional theory (DFT) using the Vienna Ab Initio Simulation Package (VASP) version 5.4.4.^{36, 37} Exchange-correlation effects were described by the functional of Perdew-Burke-Ernzerhof (PBE)³⁸ and van der Waals contributions treated with the D3 dispersion correction.³⁹ We utilized the projector augmented wave (PAW) method⁴⁰ together with an energy cutoff of 400 eV. The ionic positions of all structures were optimized within 10^{-6} eV with a Gaussian smearing of 0.01 eV. We performed two types of calculations: First for the isolated heptacene molecule and its most stable dimer⁴¹ and second for a monolayer of heptacene on Ag(110). The isolated molecules were simulated in supercells with a minimum of 15 Å vacuum between periodic replicas. For the monolayer of 7A on Ag(110), two orthogonal unit cells containing 7A oriented either along or across the close-packed Ag-rows have been considered (cf. SI p. 4ff). The unit cell with the long axis of 7A along the [001] is based on the experimental LEED structure of 5A on Ag(110),⁴² while the unit cell in which 7A lies along the $[1\bar{1}0]$ direction was adapted from the experimental LEED structure of 4A on Ag(110).⁴³ The surface is simulated within the repeated slab approach using five metallic layers and a 30 Å vacuum layer. To avoid spurious electrical fields, a dipole layer is inserted in the vacuum region.⁴⁴ The structure is optimized on a Monkhorst-Pack 5x2x1 grid of k-points constraining the coordinates of the two bottom Ag layers of the slab.⁴⁵ Subsequent to the geometry relaxation, the Kohn-Sham energies and orbitals are calculated non-self-consistently on a denser k-point mesh of 12x5x3 which is required for the simulation of the photoemission data. The angle-resolved photoemission intensity maps are calculated within the one-step model of photoemission⁴⁶ approximating the final state as a plane wave,⁴² modified by an exponential damping factor introduced between the substrate and the organic molecule in order to mimic the mean free path of the photoemitted electrons.⁴⁷

RESULTS AND DISCUSSION

Using a Knudsen type molecular evaporator, we have deposited heptacene onto the Ag(110) surface which was kept at room temperature. Figure 1a shows the momentum map of emission from a molecular state located ~ 0.8 eV below the Fermi edge, namely the HOMO of 7A, as will be unambiguously identified below. As a start, we compare the experimental to the theoretical momentum maps simulated for the HOMOs of isolated heptacene molecules as a superposition of two perpendicular orientations (Figure 1b). The simulated momentum maps of the HOMOs of a single orientation of 7A and the D_{2h} symmetrical diheptacene molecule **1** (D7A, Scheme 1) are shown in Figure 1c and 1d.

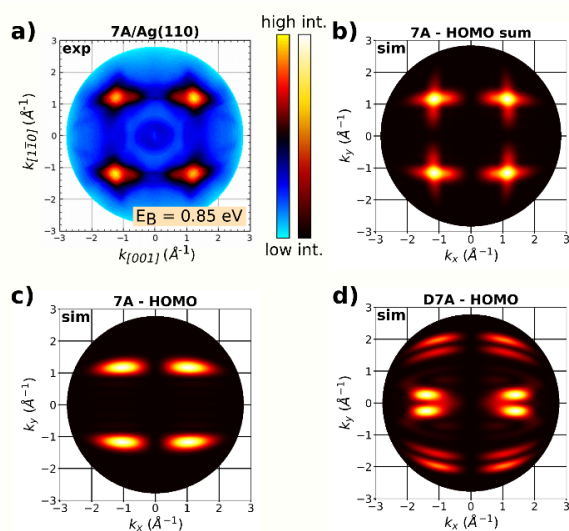


Figure 1. a) Experimental momentum map of 7A/Ag(110) measured at a binding energy of 0.85 eV, b) Simulated momentum map of the HOMO emission of two isolated heptacene molecules as a superposition of two perpendicular orientations, with a 1:1 ratio, c) Simulated momentum map

of the HOMO emission for a single heptacene orientation, d) Simulated momentum map of the HOMO emission of a single D7A molecule (1). Note that y is parallel to the long molecular axis.

The simulated HOMO map of 7A is characterized by four major lobes similar to HOMOs of other acenes (e.g., tetracene⁴⁸ and pentacene⁴²). In contrast, the map of D7A shows a splitting of the major lobes of the HOMO emission, which leads to four major lobes with much lower k_y values than for 7A and additional minor lobes (Figure 1d). Such a splitting originates from the breaking of the extended π -conjugated system into two smaller π sub-systems (one per each 7A moiety of D7A), due to the formation of C-C single bonds. However, no such splitting is observed in the experimental momentum map in Figure 1b (and the experimental maps of other molecular emissions, cf. below). Instead, the k -values of the photoemission lobes (Figure 1a) perfectly agree with those of the simulated momentum map of the 7A HOMO (Figure 1c), which suggests the formation of 7A molecules on Ag(110). Moreover, as PT is an area averaging technique, we measure the accumulated photoemission response of an ensemble of molecules on a macroscopic scale proving high reaction yield. The quality of the experimental momentum map demonstrates the chemical purity of synthesized 7A. A closer look at the experimental momentum map of Figure 1a reveals a peculiar star shaped structure of the emission pattern, instead of the elongated features expected from the simulated HOMO of 7A. Such a shape can, however, be rationalized by superimposing HOMO maps for two different orientations rotated by 90° with respect to each other. Note that the coexistence of two different adsorption configurations, i.e., orientation along both principal azimuths $[1\bar{1}0]$ and $[001]$, has been reported for shorter acenes on Ag(110) earlier.^{48, 49} Closer inspection reveals that the orientation of 7A along $[1\bar{1}0]$ and $[001]$ in a ratio of 3:1 yields the best agreement with the experiment (cf. SI Fig. S8), which is demonstrated by the

corresponding simulated map in Figure 1b. After clarifying the chemical nature and geometrical structure of the adsorbate layer, we want to analyze its electronic properties in more detail. To this end, we have recorded energy distribution maps (band maps, i.e., photoemission intensity vs (E_b , $k_{||}$)) in two different orientations of the sample ($[1\bar{1}0]$ and $[100] +45^\circ$, cf. combined band map in Figure 2a).

We have chosen these directions due to the nodal symmetry of the orbitals and molecular orientation on the surface, the LUMO and HOMO-3 of 7A have their maximal emission intensities in $[1\bar{1}0]$ -direction and the HOMO, HOMO-1 and HOMO-2 in $[100] +45^\circ$ direction. In both directions, we observe several signatures of molecular emissions marked with white symbols in Figure 2a.

Note that the π -orbitals of planar organic molecules derived from p_z -states of carbon atoms can be assigned to different π -bands of *intra*-molecular dispersion depending on their symmetry, i.e., nodal structure (cf. SI).⁵⁰ The symmetry of the orbital nodal structure in real space is reflected in reciprocal space. Thus, the π -band lowest in energy consists of orbitals concentrated at the linking and terminal carbon pairs and will further be denoted as the *linking* band. Correspondingly, the π -band next in energy consists of orbitals from the pairs of apex carbons (*apex* band) and features a characteristic node along the long molecular axis.

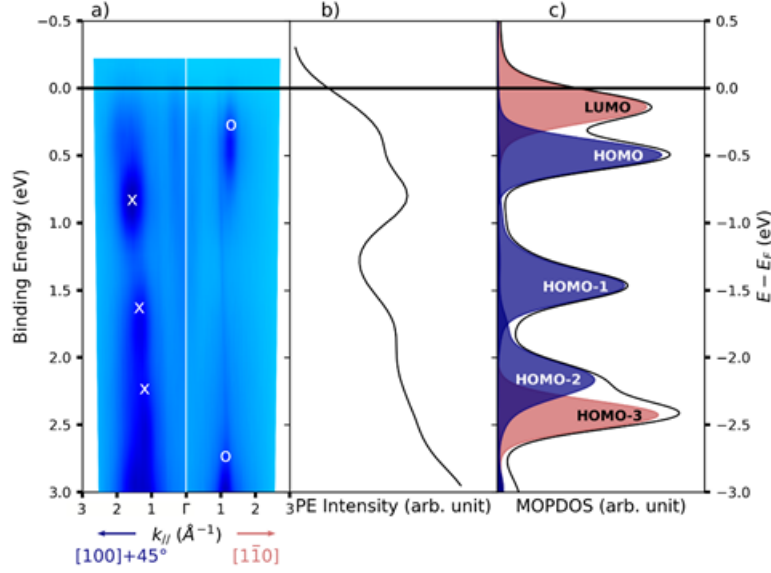


Figure 2. a) Energy distribution map (band map) in $[100]+45^\circ \rightarrow \Gamma$ and $\Gamma \rightarrow [1\bar{1}0]$ directions, white circles and crosses mark the binding energies with respect to Fermi energy at which momentum maps have been recorded. b) Energy distribution curve obtained by integration of photoemission intensity of (a) over the entire $k_{\parallel}$ range. c) MOPDOS calculated for 7A/Ag(110). Red: orbitals of the apex band, blue: orbitals of the linking band.

The spectral features of 7A in Figure 2a can accordingly be assigned to these π -bands. To do so, we, at first, compare the energy distribution curve (EDC) in Figure 2b with DFT results for a 7A monolayer on Ag(110). For comparison we have computed the density of states projected onto the molecular orbitals (MOPDOS) (Figure 2c). From the MOPDOS in Figure 2c, we also expect to observe several molecular emissions in the energy range from the Fermi level to the onset of the d-band of silver. The calculation also suggests at least partial filling of the former LUMO of 7A. Indeed, the corresponding emission feature in the experimental band map is clearly below the Fermi energy. Note that after including Hartree-Fock exchange in the exchange-correlation

functional (cf. SI), the computed molecular states shift towards higher binding energies, thereby increasing the agreement with experimental results. Emissions attributed to orbitals of the linking band are colored in red; those belonging to the apex band in blue.

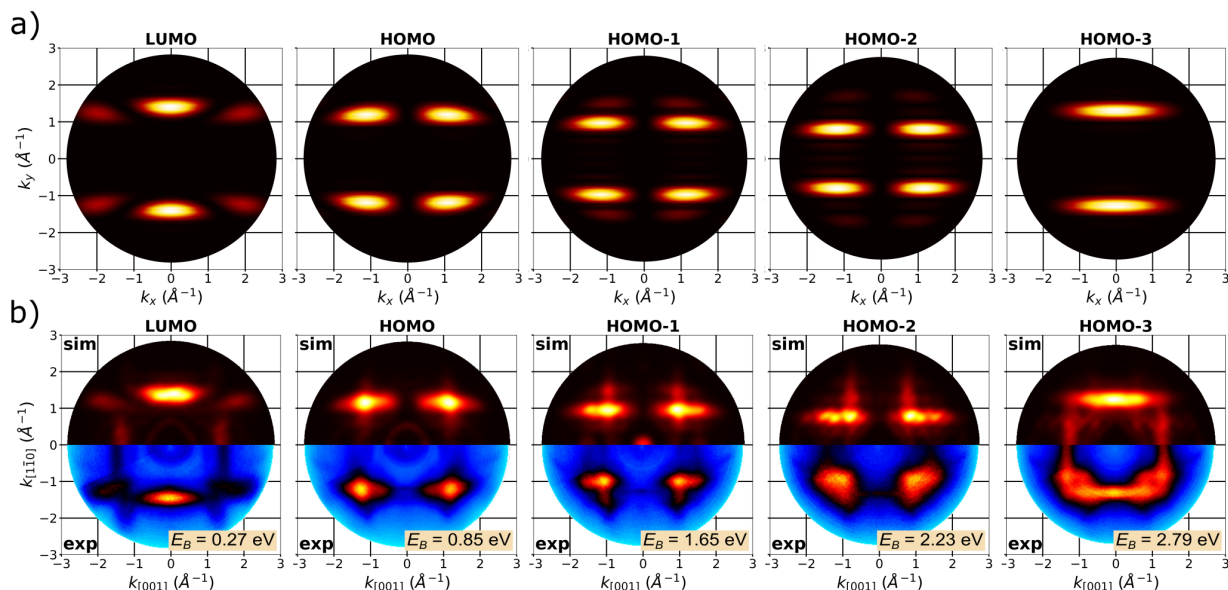


Figure 3 a) Simulated photoemission momentum maps of the π -orbitals for an isolated heptacene molecule b) Comparison of simulated (sim) and experimental (exp) photoemission momentum maps of the π -orbitals. Note that the simulated maps are superpositions of two distinct heptacene/Ag(110) azimuthal orientations in the monolayer, $[1\bar{1}0]$ and $[001]$, with a weighting ratio of 3:1.

To definitively assign the spectral features in Figure 2a,b to populated molecular orbitals of 7A, we have measured and simulated momentum maps at the binding energies marked in Figure 2a. Figure 3a shows the simulated momentum maps for an isolated heptacene molecule. In Figure 3b the experimental maps (lower halves) are compared to theoretical momentum maps simulated for the complete 7A/Ag(110) interface (upper halves) at the binding energies corresponding to the

respective peaks in the MOPDOS (Figure 2c). To mimic the experimental situation, the simulated maps are taken as a superposition of maps calculated for two heptacene orientations, along and across the silver rows in a ratio of 3:1 (cf. discussion above). The visual agreement between experimental and simulated maps allows us to unambiguously identify five molecular emissions within a small energy-window from the Fermi edge to the d-band of silver. In particular, we observe a population of the LUMO, which peaks 0.4 eV below the Fermi edge. Note that strong interactions between metal surfaces and organic molecules may change the electronic structure of the molecule at the interface resulting in charge transfer from the metal into the molecule which results in a population of the formerly unoccupied molecular states.

We note that also the profile of the C1s core level measured for 7A/Ag(110) (cf. SI) indicates a redistribution of electrons at the interface presumably due to charge transfer from metal to the molecule. At first glance, the HOMO, HOMO-1 and HOMO-2 of 7A have similar emission patterns. However, they can clearly be distinguished as their intensity maxima appear at different k -values. The HOMO has a maximum at $k_{\parallel} = 1.7 \text{ \AA}^{-1}$, HOMO-1 at $k_{\parallel} = 1.6 \text{ \AA}^{-1}$ and the maximum of HOMO-2 is at $k_{\parallel} = 1.5 \text{ \AA}^{-1}$, which is a direct consequence of the decreasing number of nodes along the long molecular axis. Following the orbitals, their measured momentum maps can also be assigned to corresponding π -bands: maps of LUMO and HOMO-3 belong to the linking band and maps of HOMO, HOMO-1 and HOMO-2 have the characteristic node along the molecular axis and thus belong to the apex band (cf. color code in Figure 2c and Figure 4).

It is important to note that, although ARPES is an area-averaging technique, we can exploit the angular dependence of the molecular emissions to shed light on differences in the electronic structure of 7A oriented parallel or perpendicular to the closed packed Ag rows, respectively. This

can be shown best for the LUMO emissions. Scanning the binding energy close to the energy of LUMO maximum and looking at the corresponding momentum maps, one notices that the partial contribution of the two perpendicularly oriented molecules varies. Namely, the LUMO of 7A molecules aligned across the Ag-rows, i.e. the $[001]$ -direction, peaks at a binding energy of about 0.1 eV lower than the LUMO of those aligned along the Ag-rows ($[1\bar{1}0]$). Deconvolution of the photoemission intensity (I , E , k_x , k_y) datacube⁴¹ using two theoretical maps of the 7A LUMO of molecules oriented perpendicular to each other confirms this finding (cf. SI), pointing at slightly different occupations of the LUMO for those two species. Note, that the energy difference derived from such deconvolution should be taken with care because of proximity to the Fermi edge. Nonetheless, we note that our DFT results for 7A molecules oriented along the $[001]$ versus the $[1\bar{1}0]$ direction on Ag(110) have also led to a binding energy shift of 0.08 eV (cf. SI), where the molecules oriented along $[001]$ have the lower energy. A shift of 0.1 eV between the two LUMOs indicates that the adsorption site (orientation) has small but not negligible influence on the electronic interaction between 7A and Ag(110). It also demonstrates the ability of photoemission tomography to draw a link between binding energies of molecules and their geometric orientation.^{33, 51-54}

When comparing our results for 7A/Ag(110) with a previous PT study of 5A/Ag(110),⁴² one can clearly recognize strong similarities, which are expected since smaller acenes develop similar orbital structures owing to their common symmetry and topology. Interestingly, the preferred orientations for 7A and 5A are different, with 7A lying predominantly along the Ag substrate rows, i.e. $[1\bar{1}0]$ orientation, while the preferred orientation of 5A molecules is across the Ag-rows i.e. along $[100]$ and rotated by 90°. This has been accounted for in Fig. 4 to facilitate comparison.

First, the LUMO of 7A adsorbed on Ag(110) appears fully occupied, while only partial LUMO occupation has been observed for 5A/Ag(110).⁴² This is expected taking into account the trend in the electron affinities, which, in the gas phase, increases from 1.01 eV for 5A to 1.57 eV for 7A, respectively.⁵⁵ Second, the energy difference between the HOMO and LUMO for 7A has shrunk to only about 0.4 eV (cf. Figure 2b) compared to 1.3 eV for 5A. This is a result of the more extended aromatic system in 7A which leads to a smaller energy spacing between electronic levels in 7A and correspondingly a reduction of the HOMO-LUMO gap. Third, increasing the length of the conjugated π -system leads to a reduction of the k-width of the features reflecting the number of scatterers in general. This can indeed be observed by comparing the LUMOs of 5A and 7A (Figure 4).

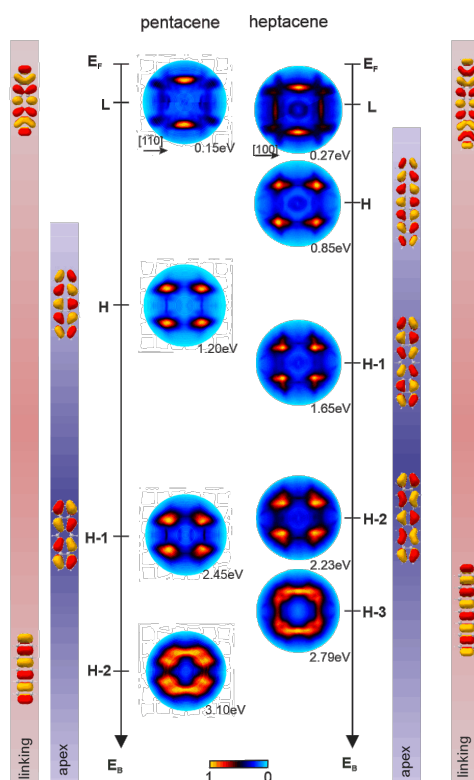


Figure 4. π -orbitals of 5A and 7A separated in two different π -bands, linking band and apex band, according to their symmetry and the corresponding experimental momentum maps of 5A/Ag(110) and 7A/Ag(110). To plot molecular orbitals, isosurfaces 10% of the maximum electron density have been used.

Finally, the overall number of π -orbitals is also increased when increasing the number of benzene units from 5A to 7A. While gas phase 5A features five occupied orbitals in the apex band and six orbitals in the linking band (cf. SI), the two additional rings of 7A give rise to four more orbitals in each band (two occupied and two unoccupied). We have measured and assigned all emissions of molecular π -states of 7A (Figure 4 right column) on Ag(110) and compared them to those of 5A (Figure 4 left column). For comparison, we also included the real space orbitals, which highlight the different nodal structures of the π -bands. The first observed pattern below the Fermi-energy belongs to the former LUMO of 5A and 7A, respectively, and can be classified into the linking band. The next molecular state (and corresponding momentum map) of the same linking band is HOMO-3 in the case of 7A and HOMO-2 in the case of 5A. Framed by these orbitals of the linking band, the topmost orbitals of the apex band are observed. Note that for 5A, PT shows the HOMO and HOMO-1 in this energy window, while for 7A three orbitals can be unambiguously distinguished: HOMO, HOMO-1 and HOMO-2 (vide supra).

CONCLUSION

To summarize, by combining ARPES experiments with ab initio electronic structure calculations, we have demonstrated the successful preparation of a well-defined monolayer of 7A on Ag(110)

via organic vapor deposition. We have utilized the power of PT to not only clearly identify intact 7A molecules and their orientations on the surface, but also to disentangle the electronic structure of the two prominent 7A adsorbate species which differ in their azimuthal alignment on the Ag(110) substrate. We find the LUMO of 7A to be occupied upon charge transfer from the Ag(110) surface and we have further characterized the electronic structure within a binding energy window from the Fermi energy down to the d-band of silver thereby identifying four more orbitals. As an outlook, succeeding in the growth of heptacene monolayers and, thus, understanding their interactions with metal contact interfaces is a prerequisite for possible applications in organic electronics. Moreover, the present results constitute a basis for future research on substituted heptacene derivatives exhibiting higher stability.

ASSOCIATED CONTENT

Supporting Information.

The Supporting Information is available free of charge on the ACS Publications website. (PDF)

- Synthesis of diheptacenes.
- Kohn-Sham orbitals assigned to the molecular π -bands of heptacene and pentacene in the gas phase.
- Calculated adsorption configurations of heptacene on Ag(110).
- Experimental EDCs compared to the density of states of heptacene/Ag(110) projected onto the molecular orbitals of a freestanding monolayer of heptacene for four different adsorption sites.
- Deconvolution of the experimental photoemission intensity $I(E_b, k_x, k_y)$ datacube with two theoretical maps of heptacene.
- C1s core level spectra of monolayers of heptacene and pentacene on Ag(110).

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Notes

The authors declare no competing financial interests.

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TOC graphic:

